# Covariant differential operators as singular intertwining operators 

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- I. Knapp-Stein intertwining operators and covariant differential operators on the sphere


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- II. Conformally invariant trilinear forms on the sphere


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- II. Conformally invariant trilinear forms on the sphere
- III. Singular invariant trilinear forms and covariant (bi-)differential operators


## Based on

- J-L. Clerc \& B. Ørsted, Conformally invariant trilinear forms on the sphere, Ann. Instit. Fourier, 61 (2011), 1807-1838


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- J-L. Clerc, T. Kobayashi, B. Ørsted \& M. Pevzner, Generalized Bernstein-Reznikov integrals, Math. Annalen 349 (2011), 395-431


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- R. Beckmann \& J-L. Clerc, Singular invariant trilinear forms and covariant (bi-)differential operators under the conformal group, J. Funct. Anal. 262 (2012), 4341-4376


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- J-L. Clerc, Singular conformally invariant trilinear forms on the sphere, in preparation

Some related references

- J. Bernstein \& A. Reznikov, Estimates of automorphic functions, Mosc. Math. J. 4 (2004), 19-37

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- T. Kobayashi and B. Speh, Symmetry breaking for representations of rank one orthogonal groups, posted on arXiv (2013)
- P. Somberg, Rankin-Cohen brackets for orthogonal Lie algebras and bilinear conformally invariant differential operators, posted on arXiv (2013)

Knapp-Stein intertwining operators and covariant differential operators on the sphere

Jean-Louis Clerc

## The geometric context

- $S=S^{n-1}$ the unit sphere in $E=\mathbb{R}^{n}$
$\mathbf{1}=(1,0, \ldots, 0)$
$S \simeq K / M$, where $K=S O(n), M \simeq S O(n-1)$


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- $S=S^{n-1}$ the unit sphere in $E=\mathbb{R}^{n}$
$\mathbf{1}=(1,0, \ldots, 0)$
$S \simeq K / M$, where $K=S O(n), M \simeq S O(n-1)$
- $\mathcal{S}=$ set of isotropic lines in the Lorentz space $E^{1, n}$.

$$
S \ni x \longmapsto \mathbb{R}(1, x) \in \mathcal{S}
$$

is a 1-1 correspondence, and so $G$ acts on $S$.

- The Lorentz group $G=S O_{0}(1, n)$ acts on $\mathcal{S}$, hence on $S$. $S \simeq G / P$, where $P$ the stabilizer of $\mathbf{1}$ in $G$ is a parabolic subgroup of $G$.
- The action of $G$ is conformal, i.e. for any tangent vector $\xi$ to $S$ at $x$

$$
|D g(x) \xi|=\kappa(g, x)|\xi|
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where $\kappa(g, x)$ is the conformal factor of $g$ at $x$.

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- Covariance property of the Euclidean distance on $S$ :

$$
|g(x)-g(y)|=\kappa(g, x)^{\frac{1}{2}}|x-y| \kappa(g, y)^{\frac{1}{2}}
$$

for $x, y \in S$ and $g \in G$.

## The principal series

- The group $G$ acts naturally on the spaces of smooth densities on $S$, thus defining the (scalar) principal series of representations of $G$.
- Let $d x$ be the Lebesgue measure on the sphere, use it to identify smoth densities with $\mathcal{C}^{\infty}$ functions. For $\lambda \in \mathbb{C}$, we obtain the following representation of $G$ on $\mathcal{C}^{\infty}(S)$

$$
\pi_{\lambda}(g) f(x)=\kappa\left(g^{-1}, x\right)^{\frac{n-1}{2}+\lambda} f\left(g^{-1}(x)\right)
$$

[The shift by $\rho=\frac{n-1}{2}$ is to have $\pi_{\lambda}$ unitary for pure imaginary $\lambda$ ]

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- The duality relation between $\pi_{\lambda}$ and $\pi_{-\lambda}$

$$
\int_{S} \pi_{\lambda}(g) \varphi(x) \psi(x) d x=\int_{S} \varphi(x) \pi_{-\lambda}\left(g^{-1}\right) \psi(x) d x
$$

## The Knapp-Stein intertwining operators

- Knapp-Stein operator (formal)

$$
J_{\lambda} f(x)=\int_{S}|x-y|^{-(n-1)+2 \lambda} f(y) d y
$$

- intertwining relation

$$
J_{\lambda} \circ \pi_{\lambda}(g)=\pi_{-\lambda}(g) \circ J_{\lambda}, \quad \text { for any } g \in G
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- convergent for $\Re \lambda>0$
- look for analytic continuation of $J_{\lambda}$ for $\lambda \in \mathbb{C}$


## Bernstein-Sato identity

For $f$ a given real analytic function everywhere $\geq 0$

$$
D\left(x, \frac{\partial}{\partial x}, s\right) f^{s+1}(x)=b(s) f^{s}(x)
$$

valid on $\{x, f(x) \neq 0\}$.

- $D\left(x, \frac{\partial}{\partial x}, s\right)$ is a differential operator with smooth coefficients, polynomial in the parameter $s \in \mathbb{C}$.
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- example : BS identity for $|\mathbf{1}-x|^{2}$ on $S$

$$
\left(\Delta+\left(\frac{s}{2}+1\right)\left(\frac{s}{2}+n-1\right)\right)|\mathbf{1}-x|^{s+2}=(s+2)(s+n-1)|\mathbf{1}-x|^{s}
$$

where $\Delta$ is the Laplacian on $S$.

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where $\Delta$ is the Laplacian on $S$.
Hint: $|\mathbf{1}-x|^{2}=2\left(1-x_{1}\right)$ and

$$
\Delta \varphi\left(x_{1}\right)=\left(1-x_{1}^{2}\right) \varphi^{\prime \prime}\left(x_{1}\right)-(n-1) \varphi^{\prime}\left(x_{1}\right) .
$$

## Meromorphic continuation of $\int_{S}|x-y|^{S} f(y) d y$

- from the previous Bernstein-Sato identity

$$
|x-y|^{s}=\frac{1}{(s+2)(s+n-1)}\left(\Delta_{y}+\left(\frac{s}{2}+1\right)\left(\frac{s}{2}+n-1\right)\right)|x-y|^{s+2}
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- integrate by parts to get

$$
\begin{gathered}
\int_{S}|x-y|^{s} f(y) d y= \\
\frac{1}{(s+2)(s+n-1)} \int_{S}|x-y|^{s+2}\left(\Delta_{y}+\left(\frac{s}{2}+1\right)\left(\frac{s}{2}+n-1\right) f\right)(y) d y
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\end{gathered}
$$

- this allows the meromorphic continuation, poles occurring for $s=-(n-1),-(n-1)-2, \ldots,-(n-1)-2 k, \ldots$


## Meromorphic continuation of $J_{\lambda}$

- Letting $s=-(n-1)+2 \lambda$, the operator-valued function $\lambda \rightarrow J_{\lambda}$ can be continued meromorphically to $\mathbb{C}$ with poles at $-k, k \in \mathbb{N}$.


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- for $\lambda \sim 0$, write $J_{\lambda} f(x)=\int_{S}|x-y|^{-(n-1)+2 \lambda} f(y) d y$ as

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\int_{S}|x-y|^{-(n-1)+2 \lambda}(f(y)-f(x)) d y+f(x) \int_{S}|x-y|^{-(n-1)+2 \lambda} d y
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- $f(x)-f(y)=O(|x-y|)$, hence no problem with the first integral

$$
\int_{S}|x-y|^{s} d y=\int_{S}|\mathbf{1}-y|^{s} d y=(2 \sqrt{\pi})^{n-1} 2^{\frac{\Gamma\left(\frac{s}{2}+\rho\right)}{\Gamma\left(\frac{s}{2}+2 \rho\right)}}
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$$

- hence

$$
\operatorname{Res}\left(J_{\lambda} f(x), \lambda=0\right)=\frac{\pi^{\rho}}{\Gamma(\rho)} f(x)
$$

- for $\lambda \sim-1$, use BS identity, integrate by parts, take residues at -1 on both sides to get

$$
\operatorname{Res}\left(J_{\lambda} f(x), \lambda=-1\right)=\frac{\pi^{\rho} 4^{-1}}{\Gamma(\rho+1)} \Delta_{1} f(x)
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where $\Delta_{1}=\Delta-\frac{1}{4}(n-1)(n-3)$ is the conformal Laplacian (Yamabe operator) on $S$.

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- iterating the argument

$$
\operatorname{Res}\left(J_{\lambda}, \lambda=-k\right)=\frac{\pi^{\rho} 4^{-k}}{\Gamma(\rho+k)} \Delta_{k}
$$

where

$$
\Delta_{k}=\prod_{j=1}^{k}(\Delta-(\rho+j-1)(\rho-j))
$$

## Covariant differential operators for densities on $S$

The intertwining property satisfied by $J_{\lambda}$ between $\pi_{\lambda}, \pi_{-\lambda}$ extends to the residues and yields

$$
\Delta_{k} \circ \pi_{-k}(g)=\pi_{k}(g) \circ \Delta_{k}
$$

usually know as the covariance property for $\Delta_{k}$.

It can be shown that there are no more covariant differential operators between spaces of densities. The proof uses results on homomorphisms of (generalized) Verma modules.

## Sketch of a down-to-earth proof

- Use the noncompact picture $\mathbb{R}^{n-1}$.
- Let $D$ be a differential operator on $\mathbb{R}^{n-1}$ covariant w.r.t. $\left(\pi_{\lambda}, \pi_{\mu}\right)$.


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- $A$ acts by dilation on $\mathbb{R}^{n-1}$, hence $D$ has to be homogeneous. So $D=a_{k} \Delta^{k}$ for some $k \in \mathbb{N}$, and $\mu=\lambda+2 k$.


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- $A$ acts by dilation on $\mathbb{R}^{n-1}$, hence $D$ has to be homogeneous. So $D=a_{k} \Delta^{k}$ for some $k \in \mathbb{N}$, and $\mu=\lambda+2 k$.
- If $D$ is non trivial, then the Casimir operator acts on both spaces by the same scalar. This forces $\lambda=-\mu=-k$. Q.E.D.


## End of part I

Thank you for your attention!

