# G-STRUCTURES WITH PRESCRIBED GEOMETRY—EXAMPLES 

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References: For these three lectures (these are the slide for the third), the reader who wants more background information on exterior differential systems might want to consult the brief introduction
http://www.math.duke.edu/~bryant/Introduction_to_EDS.pdf

Many of the examples discussed here and the main variants of Cartan's theory of structure equations can be found in the lecture notes on EDS that can be found here

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http://arxiv.org/abs/1405.3116
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This latter article contains many references to the literature and further resources.

## I. Riemannian Surfaces with $|\nabla K|^{2}=1$

Consider $\left(M^{2}, g\right)$ whose Gauss curvature satisfies $|\nabla K|^{2}=1$. The structure equations on the orthonormal frame bundle $\pi: B \rightarrow M$ have $g=\omega_{1}{ }^{2}+\omega_{2}{ }^{2}$ and are

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\begin{aligned}
\mathrm{d} \omega_{1} & =-\omega_{12} \wedge \omega_{2} \\
\mathrm{~d} \omega_{2} & =\omega_{12} \wedge \omega_{1} \\
\mathrm{~d} \omega_{12} & =a \omega_{1} \wedge \omega_{2} \\
\mathrm{~d} a & =\cos b \omega_{1}+\sin b \omega_{2}
\end{aligned} \quad \omega_{1} \wedge \omega_{2} \wedge \omega_{12} \neq 0
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where $K=a$ is the Gauss curvature and $b$ is the free derivative.

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where $K=a$ is the Gauss curvature and $b$ is the free derivative.
Using these equations, we see that $\mathrm{d}^{2} \omega_{1}=\mathrm{d}^{2} \omega_{2}=\mathrm{d}^{2} \omega_{12}=0$ are identities, but, using the structure equations, one finds

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0=\mathrm{d}(\mathrm{~d} a)=\left(\mathrm{d} b-\omega_{12}\right) \wedge\left(-\sin b \omega_{1}+\cos b \omega_{2}\right)
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It follows that the hypotheses of the Variant of Cartan's Theorem are satisfied, with the characters of the tableau of free derivatives being $s_{1}=1$, $s_{2}=s_{3}=0$. Thus, the general (local) solution depends on one function of one variable.
I. Ricci solitons in dimension 2

Consider $\left(M^{2}, g\right)$ that is a gradient Ricci soliton, i.e.,

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\operatorname{Ric}(g)=K g=\operatorname{Hess}(f)=\nabla^{2} f
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for some function $f$ on $M$.

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\mathrm{~d} \omega_{12} & =K \omega_{1} \wedge \omega_{2}
\end{aligned}
$$

$$
\begin{aligned}
\mathrm{d} K & =K_{1} \omega_{1}+K_{2} \omega_{2} \\
\mathrm{~d} f & =f_{1} \omega_{1}+f_{2} \omega_{2} \\
\mathrm{~d} f_{1} & =-f_{2} \omega_{12}+K \omega_{1} \\
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\mathrm{d} \omega_{1} & =-\omega_{12} \wedge \omega_{2} & \mathrm{~d} K & =K_{1} \omega_{1}+K_{2} \omega_{2} \\
\mathrm{~d} \omega_{2}= & \omega_{12} \wedge \omega_{1} & \mathrm{~d} f & =f_{1} \omega_{1}+f_{2} \omega_{2} \\
\mathrm{~d} \omega_{12} & =K \omega_{1} \wedge \omega_{2} & \mathrm{~d} f_{1} & =-f_{2} \omega_{12}+K \omega_{1} \\
\mathrm{~d} f_{2} & =f_{1} \omega_{12}+K \omega_{2} \\
& 0=\mathrm{d}\left(\mathrm{~d} f_{1}\right)=\left(K_{2}+K f_{2}\right) \omega_{1} \wedge \omega_{2} . \\
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\mathrm{~d} \omega_{12} & =K \omega_{1} \wedge \omega_{2}
\end{aligned}
$$

$$
\mathrm{d} K=-K\left(f_{1} \omega_{1}+f_{2} \omega_{2}\right)
$$

$$
\mathrm{d} f=f_{1} \omega_{1}+f_{2} \omega_{2}
$$

$$
\mathrm{d} f_{1}=-f_{2} \omega_{12}+K \omega_{1}
$$

$$
\mathrm{d} f_{2}=f_{1} \omega_{12}+K \omega_{2}
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& 0=\mathrm{d}\left(\mathrm{~d} f_{1}\right)=\left(K_{2}+K f_{2}\right) \omega_{1} \wedge \omega_{2} \\
& 0=\mathrm{d}\left(\mathrm{~d} f_{2}\right)=\left(K_{1}+K f_{1}\right) \omega_{2} \wedge \omega_{1}
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Thus, the above structure equations can be tightened to

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\mathrm{d} K=-K\left(f_{1} \omega_{1}+f_{2} \omega_{2}\right)=-K \mathrm{~d} f .
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\begin{array}{rlrl}
\mathrm{d} \omega_{1}=-\omega_{12} & \wedge \omega_{2} & \mathrm{~d} K & =-K\left(f_{1} \omega_{1}+f_{2} \omega_{2}\right) \\
\mathrm{d} \omega_{2}=\omega_{12} & \wedge \omega_{1} & \mathrm{~d} f & =f_{1} \omega_{1}+f_{2} \omega_{2} \\
\mathrm{~d} \omega_{12}=K \omega_{1} & \wedge \omega_{2} & \mathrm{~d} f_{1} & =-f_{2} \omega_{12}+K \omega_{1} \\
\mathrm{~d} f_{2} & =f_{1} \omega_{12}+K \omega_{2} \\
0=\mathrm{d}\left(\mathrm{~d} f_{1}\right)=\left(K_{2}+K f_{2}\right) \omega_{1} \wedge \omega_{2} \\
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There are now no more 'free derivatives', but $\mathrm{d}^{2}=0$ is an identity.
III. Prescribed curvature equations for Finsler surfaces

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For an oriented Finsler surface ( $\left.M^{2}, F\right)$, Cartan showed that the 'tangent indicatrix' $\Sigma \subset T M$ has a canonical coframing ( $\omega_{1}, \omega_{2}, \omega_{3}$ ), satisfying

$$
\begin{array}{ll}
\mathrm{d} \omega_{1}=-\omega_{2} \wedge \omega_{3} \\
\mathrm{~d} \omega_{2}=-\omega_{3} \wedge \omega_{1}-I \omega_{2} \wedge \omega_{3} \\
\mathrm{~d} \omega_{3}=-K \omega_{1} \wedge \omega_{2}-J \omega_{2} \wedge \omega_{3} & \omega_{1} \wedge \omega_{2} \wedge \omega_{3} \neq 0 \\
\hline
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where I have written $\omega_{3}$ for what would be $-\omega_{12}$ in the Riemannian case.
The functions $I, J$, and $K$ are the Finsler structure functions. $K$ is the Finsler-Gauss (or 'flag') curvature; $I$ is the Cartan scalar, which vanishes iff $\left(M^{2}, F\right)$ is Riemannian; $J$ is the Landsberg scalar.

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Differentiating the above equations yields the Finsler-Bianchi identities

$$
\begin{array}{rrrr}
\mathrm{d} I= & J \omega_{1} & +I_{2} \omega_{2} & +I_{3} \omega_{3}, \\
\mathrm{~d} J= & -\left(K_{3}+K I\right) \omega_{1} & +J_{2} \omega_{2} & +J_{3} \omega_{3}, \\
\mathrm{~d} K= & K_{1} \omega_{1} & +K_{2} \omega_{2} & +K_{3} \omega_{3} .
\end{array}
$$

for seven new functions, $K_{i}$, etc. These are the 'free derivatives', and their tableau has $\left(s_{1}, s_{2}, s_{3}\right)=(3,3,1)$.

Case 1: $I=0$ (the Riemannian case), forces $J=0$, so

$$
\begin{aligned}
\mathrm{d} \omega_{1} & =-\omega_{2} \wedge \omega_{3} \\
\mathrm{~d} \omega_{2} & =-\omega_{3} \wedge \omega_{1} \\
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\mathrm{~d} K & =K_{1} \omega_{1}+K_{2} \omega_{2}
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Case 2: $J=0$ (Landsberg surfaces)

$$
\begin{array}{rlrr}
\mathrm{d} I & = & & +I_{2} \omega_{2} \\
\mathrm{~d} K & = & +I_{3} \omega_{3}, \\
& K_{1} \omega_{1} & +K_{2} \omega_{2} & -K I \omega_{3} .
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& K_{1} \omega_{1} & +K_{2} \omega_{2} & -K I \omega_{3} .
\end{array}
$$

Tableau has $\left(s_{1}, s_{2}, s_{3}\right)=(2,2,0)$, so 2 functions of 2 variables.

Case 3: $K$-basic (i.e., $K_{3}=0$, so $K$ well-defined on $M$ )

$$
\begin{array}{rrrrr}
\mathrm{d} I= & J \omega_{1} & +I_{2} \omega_{2} & +I_{3} \omega_{3}, \\
\mathrm{~d} J & = & -K I \omega_{1} & +J_{2} \omega_{2} & +J_{3} \omega_{3}, \\
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Case 4: $K$ is constant

$$
\begin{array}{rrrrr}
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\mathrm{~d} J & = & -K I \omega_{1} & +J_{2} \omega_{2} & +J_{3} \omega_{3}, \\
\mathrm{~d} K & = & 0 &
\end{array}
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## IV. Ricci potentials in dimension 3

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Consider the problem of the generality of Riemannian 3-manifolds $\left(M^{3}, g\right)$ for which there exists a function $f$ (a 'Ricci potential') such that

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\operatorname{Ric}(g)=(\mathrm{d} f)^{2}+H(f) g
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where $H$ is a specified function of one variable. On the orthonormal frame bundle $B^{6} \rightarrow M^{3}$, we have structure equations

$$
\mathrm{d} \omega_{i}=-\omega_{i j} \wedge \omega_{j} \quad \text { and } \quad\left(\begin{array}{l}
\mathrm{d} \omega_{23} \\
\mathrm{~d} \omega_{31} \\
\mathrm{~d} \omega_{12}
\end{array}\right)=-\left(\begin{array}{l}
\omega_{12} \wedge \omega_{31} \\
\omega_{23} \wedge \omega_{12} \\
\omega_{31} \wedge \omega_{23}
\end{array}\right)-\left(R-\frac{1}{2} \operatorname{tr}(R) I_{3}\right)\left(\begin{array}{l}
\omega_{2} \wedge \omega_{3} \\
\omega_{3} \wedge \omega_{1} \\
\omega_{1} \wedge \omega_{2}
\end{array}\right)
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where $R=\left(R_{i j}\right)$ is the symmetric matrix of the Ricci tensor.

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\end{array}\right)
$$

where $R=\left(R_{i j}\right)$ is the symmetric matrix of the Ricci tensor. By hypothesis, there exists a function $f$ such that $R_{i j}=f_{i} f_{j}+H(f) \delta_{i j}$ where

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\mathrm{d} f=f_{1} \omega_{1}+f_{2} \omega_{2}+f_{3} \omega_{3}
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\mathrm{d} f=f_{1} \omega_{1}+f_{2} \omega_{2}+f_{3} \omega_{3} .
$$

The four functions $\left(f, f_{1}, f_{2}, f_{3}\right)$ will play the role of the $a^{\alpha}$ in the structure equations, and $\mathrm{d}(\mathrm{d} f)=0$ implies that there exist $f_{i j}=f_{j i}$ so that

$$
\mathrm{d} f_{i}=-\omega_{i j} f_{j}+f_{i j} \omega_{j}
$$

The equations $\mathrm{d}\left(\mathrm{d} \omega_{i}\right)=0$ are identities (because $R$ is symmetric), but the equations $\mathrm{d}\left(\mathrm{d} \omega_{i j}\right)=0$ can be written as

$$
\left(2\left(f_{11}+f_{22}+f_{33}\right)-H^{\prime}(f)\right) \mathrm{d} f=0 .
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$$
\left(2\left(f_{11}+f_{22}+f_{33}\right)-H^{\prime}(f)\right) \mathrm{d} f=0
$$

Thus, either $\mathrm{d} f=0$ (in which case, the metric is Einstein) or else

$$
f_{11}+f_{22}+f_{33}-\frac{1}{2} H^{\prime}(f)=0
$$

so that one has

$$
\mathrm{d} f_{i}=-\omega_{i j} f_{j}+\left(b_{i j}+\frac{1}{6} H^{\prime}(f) \delta_{i j}\right) \omega_{j} .
$$

where the (new) $b_{i j}=b_{j i}$ are subject to the trace condition $b_{11}+b_{22}+b_{33}=0$. These $b_{i j}$ will play the role of the 'free derivatives' in the structure equations.

The equations $\mathrm{d}\left(\mathrm{d} \omega_{i}\right)=0$ are identities (because $R$ is symmetric), but the equations $\mathrm{d}\left(\mathrm{d} \omega_{i j}\right)=0$ can be written as

$$
\left(2\left(f_{11}+f_{22}+f_{33}\right)-H^{\prime}(f)\right) \mathrm{d} f=0
$$

Thus, either $\mathrm{d} f=0$ (in which case, the metric is Einstein) or else

$$
f_{11}+f_{22}+f_{33}-\frac{1}{2} H^{\prime}(f)=0,
$$

so that one has

$$
\mathrm{d} f_{i}=-\omega_{i j} f_{j}+\left(b_{i j}+\frac{1}{6} H^{\prime}(f) \delta_{i j}\right) \omega_{j} .
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where the (new) $b_{i j}=b_{j i}$ are subject to the trace condition $b_{11}+b_{22}+b_{33}=0$. These $b_{i j}$ will play the role of the 'free derivatives' in the structure equations.

We can now easily check that these structure equations satisfy the compatibility and involutivity conditions for the Cartan Variant Theorem, with characters

$$
\left(s_{1}, s_{2}, s_{3}, s_{4}, s_{5}, s_{6}\right)=(3,2,0,0,0,0)
$$

so, up to diffeomorphism, these structures $\left(M^{3}, g, f\right)$ depend on 2 functions of 2 variables.
V. Einstein-Weyl structures in dimension 3

## V. Einstein-Weyl structures in dimension 3

These are $\mathrm{CO}(3)$-structures on 3 -manifolds endowed with a compatible connection on $B^{7} \rightarrow M^{3}$ that satisfies the structure equations

$$
\left(\begin{array}{l}
\mathrm{d} \eta_{1} \\
\mathrm{~d} \eta_{2} \\
\mathrm{~d} \eta_{3}
\end{array}\right)=-\left(\begin{array}{rrr}
\theta_{0} & \theta_{3} & -\theta_{2} \\
-\theta_{3} & \theta_{0} & \theta_{1} \\
\theta_{2} & -\theta_{1} & \theta_{0}
\end{array}\right) \wedge\left(\begin{array}{l}
\eta_{1} \\
\eta_{2} \\
\eta_{3}
\end{array}\right)
$$

and

$$
\left(\begin{array}{c}
\mathrm{d} \theta_{0} \\
\mathrm{~d} \theta_{1} \\
\mathrm{~d} \theta_{2} \\
\mathrm{~d} \theta_{3}
\end{array}\right)=\left(\begin{array}{c}
0 \\
\theta_{2} \wedge \theta_{3} \\
\theta_{3} \wedge \theta_{1} \\
\theta_{1} \wedge \theta_{2}
\end{array}\right)+\left(\begin{array}{ccc}
2 H_{1} & 2 H_{2} & 2 H_{3} \\
H_{0} & H_{3} & -H_{2} \\
-H_{3} & H_{0} & H_{1} \\
H_{2} & -H_{1} & H_{0}
\end{array}\right)\left(\begin{array}{l}
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\end{array}\right) .
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\end{array}\right)\left(\begin{array}{l}
\eta_{2} \wedge \eta_{3} \\
\eta_{3} \wedge \eta_{1} \\
\eta_{1} \wedge \eta_{2}
\end{array}\right) .
$$

Differentiating these equations gives relations of the following form ( $(i, j, k)$ is an even permutation of $(1,2,3))$, where, the $9 L_{i j}$ satisfy $L_{11}+L_{22}+L_{33}=0$.

$$
\begin{aligned}
\mathrm{d} H_{0} & =2 H_{0} \theta_{0}+\left(L_{i j}-L_{j i}\right) \eta_{k} \\
\mathrm{~d} H_{i} & =2 H_{i} \theta_{0}-H_{j} \theta_{k}+H_{k} \theta_{j}+L_{i j} \eta_{j}
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H_{0} & H_{3} & -H_{2} \\
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$$

Differentiating these equations gives relations of the following form ( $(\mathrm{i}, \mathrm{j}, \mathrm{k})$ is an even permutation of $(1,2,3)$ ), where, the $9 L_{i j}$ satisfy $L_{11}+L_{22}+L_{33}=0$.

$$
\begin{aligned}
\mathrm{d} H_{0} & =2 H_{0} \theta_{0}+\left(L_{i j}-L_{j i}\right) \eta_{k} \\
\mathrm{~d} H_{i} & =2 H_{i} \theta_{0}-H_{j} \theta_{k}+H_{k} \theta_{j}+L_{i j} \eta_{j}
\end{aligned}
$$

These are involutive with $\left(s_{1}, s_{2}, s_{3}, \ldots\right)=(4,4,0, \ldots)$.
VI. Nearly-Kähler 6-manifolds

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A nearly-Kähler 6 -manifold is a manifold $M^{6}$ endowed with a $\mathrm{SU}(3)$-structure $\pi: B \rightarrow M$ with torsion, whose first structure equations take the form

$$
\mathrm{d} \eta_{i}=-\theta_{i \bar{l}} \wedge \eta_{l}+\lambda \overline{\eta_{j}} \wedge \overline{\eta_{k}}
$$

where $(i, j, k)$ is an even perm. of $(1,2,3), \overline{\theta_{i \bar{\jmath}}}=-\theta_{j \bar{\imath}}$ and $\theta_{i \bar{\imath}}=0$, and $\lambda=\bar{\lambda}$.

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The identities $\mathrm{d}\left(\mathrm{d} \eta_{i}\right)=0$ imply that $\mathrm{d} \lambda=0$ and

$$
\mathrm{d} \theta_{i \bar{\jmath}}=-\theta_{i \bar{k}} \wedge \theta_{k \bar{\jmath}}+\lambda^{2}\left(\frac{3}{4} \eta_{i} \wedge \overline{\eta_{j}}-\frac{1}{4} \delta_{i \bar{\jmath}} \eta_{l} \wedge \overline{\eta_{l}}\right)+K_{i \bar{\jmath} p \bar{q}} \eta_{q} \wedge \overline{\eta_{p}},
$$

where the functions $K_{i \bar{\jmath} p \bar{q}}$ satisfy

$$
K_{i \bar{\jmath} p \bar{q}}=K_{p \bar{j} \bar{q} \bar{q}}=K_{i \bar{q} p \bar{\jmath}}=\overline{K_{j \bar{\imath} q \bar{p}}} \quad \text { and } \quad K_{i \bar{\imath} p \bar{q}}=0 .
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$$

This leaves 27 real components in the tensor $K$, and $\mathrm{d}\left(\mathrm{d} \theta_{i \bar{\jmath}}\right)=0$ yields
$\mathrm{d} K_{i \bar{\jmath} p \bar{q}}=-K_{\ell \bar{\jmath} p \bar{q}} \theta_{i \bar{\ell}}+K_{i \bar{\ell} \bar{q} \bar{q}} \theta_{\ell \bar{\jmath}}-K_{i \bar{j} \ell \bar{q}} \theta_{p \bar{\ell}}+K_{i \bar{\jmath} p \bar{\ell}} \theta_{\ell \bar{q}}+K_{i \bar{\jmath} \bar{q} \bar{\ell} \bar{\ell}} \eta_{\ell}+K_{i \bar{\jmath} p \bar{q} \ell} \eta_{\bar{\ell}}$, where $K_{i \bar{\jmath} p \bar{q} \bar{\ell}}=K_{i \bar{\jmath} p \bar{\ell} \bar{q}}$ and $K_{i \bar{\jmath} p \bar{q} \ell}=K_{i \bar{\jmath} \bar{q} \bar{q}}$, etc.

## VI. Nearly-Kähler 6-manifolds

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$$
\mathrm{d} \eta_{i}=-\theta_{i \bar{l}} \wedge \eta_{l}+\lambda \overline{\eta_{j}} \wedge \overline{\eta_{k}}
$$

where $(i, j, k)$ is an even perm. of $(1,2,3), \overline{\theta_{i \bar{\jmath}}}=-\theta_{j \bar{\imath}}$ and $\theta_{i \bar{\imath}}=0$, and $\lambda=\bar{\lambda}$.
The identities $\mathrm{d}\left(\mathrm{d} \eta_{i}\right)=0$ imply that $\mathrm{d} \lambda=0$ and

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$$

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K_{i \bar{\jmath} p \bar{q}}=K_{p \bar{j} \bar{q} \bar{q}}=K_{i \bar{q} p \bar{\jmath}}=\overline{K_{j \bar{\imath} q \bar{p}}} \quad \text { and } \quad K_{i \bar{\imath} p \bar{q}}=0 .
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This is involutive, with characters $\left(s_{1}, \ldots, s_{5}, s_{6}\right)=(27,27,19,9,2,0)$.

