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New D-term uplift

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How is N=1 SUSY broken?

- F-term (chiral multiplet).
- **D-term** ($U(1)$ gauge multiplet).
- Other methods (complex linear, etc).

Fayet–Iliopoulos D-term

→ A gauge multiplet contains the component fields

$$V_{WZ} = -\theta\sigma^m\bar{\theta}A_m - i\bar{\theta}^2\theta^\alpha\lambda_\alpha + i\theta^2\bar{\theta}_{\dot{\alpha}}\bar{\lambda}^{\dot{\alpha}} + \frac{1}{2}\theta^2\bar{\theta}^2D.$$

→ The simplest model with a Fayet–Iliopoulos term is

$$\begin{aligned}\mathcal{L} &= \frac{1}{4} \left(\int d^2\theta W^2(V) + c.c. \right) - 2\xi \int d^4\theta V \\ &= -\frac{1}{4} F^{mn}F_{mn} - i\lambda\sigma^m\partial_m\bar{\lambda} + \frac{1}{2}D^2 - \xi D.\end{aligned}$$

Fayet, Iliopoulos '74

Supersymmetry is broken spontaneously

→ The auxiliary field gets a vev: $\langle D \rangle = \xi$.

→ The goldstino is: $\delta\lambda_\beta = -i\xi\epsilon_\beta + \dots$

→ The vacuum energy is: $\langle V \rangle = \frac{1}{2}\xi^2$.

How do we couple to supergravity?

Plan

- Freedman model and R-symmetry gauging.
- New model without R-symmetry gauging.
- Relation to KKLT-type uplift.

The Freedman model

The Noether method

- ▶ The gauge field of SUSY is the gravitino ψ_m , with supersymmetry transformation

$$\delta\psi_{m\alpha} = -2\partial_m\epsilon_\alpha + \dots$$

- ▶ We start from $-e\xi D$ and perform the Noether procedure

$$\partial_m\epsilon_\alpha(x) \rightarrow \psi_{m\alpha}.$$

- ▶ At some point we have to cancel

$$-ie\xi\epsilon^{klmn}(\bar{\psi}_k\bar{\sigma}_l\epsilon - \bar{\epsilon}\bar{\sigma}_l\psi_k)\partial_n A_m.$$

- ▶ Therefore we have to add

$$\mathcal{L}_{\psi\psi A} = \frac{i}{2} e \xi \epsilon^{klmn} \bar{\psi}_k \bar{\sigma}_l \psi_n A_m.$$

- ▶ This means that the gravitino is charged under the $U(1)$.

→ The FI term in supergravity requires R-symmetry gauging by A_m .

→ This has impact on model building. (A gravitino mass is in some cases not allowed.) *Barbieri, Ferrara, Nanopoulos, Stelle '82, Villadoro, Zwirner '05*

The complete model

Freedman '77

- ▶ Once we complete the procedure we have (on-shell)

$$e^{-1} \mathcal{L} \Big|_{\lambda=0} = -\frac{1}{2} R + \frac{1}{2} \epsilon^{klmn} (\bar{\psi}_k \bar{\sigma}_l \mathcal{D}_m \psi_n - \psi_k \sigma_l \mathcal{D}_m \bar{\psi}_n) \\ - \frac{1}{4} F_{mn} F^{mn} + \frac{i}{2} e \xi \epsilon^{klmn} \bar{\psi}_k \bar{\sigma}_l \psi_n A_m - \frac{1}{2} \xi^2 .$$

- ▶ In superspace this would read

$$\mathcal{L} = -3 \int d^4 \theta E e^{\frac{2}{3} \xi V} + \frac{1}{4} \left(\int d^2 \Theta 2 \mathcal{E} W^2(V) + c.c. \right) .$$

A new supergravity embedding
Cribiori, Tournoy, FF, Van Proeyen '17

- ▶ We observe that λ_α/D has a very nice property

$$\delta \left(i \frac{\lambda_\alpha}{D} \right) = \epsilon_\alpha + \dots$$

- ▶ To cancel the term

$$-ie\xi \epsilon^{klmn} (\bar{\psi}_k \bar{\sigma}_l \epsilon - \bar{\epsilon} \bar{\sigma}_l \psi_k) \partial_n \mathbf{A}_m$$

during the Noether procedure, we introduce instead the term

$$-e\xi \epsilon^{klmn} \left(\bar{\psi}_k \bar{\sigma}_l \frac{\lambda}{D} - \frac{\bar{\lambda}}{D} \bar{\sigma}_l \psi_k \right) \partial_n \mathbf{A}_m.$$

Supersymmetry has to be broken and by assumption: $\langle D \rangle \neq 0$.

A new FI D-term

- ▶ In superspace the full $U(1)$ sector is

$$\begin{aligned}\mathcal{L}_{NEW} = & \frac{1}{4} \left(\int d^2\Theta 2\mathcal{E} W^2(V) + c.c. \right) \\ & + 8\xi \int d^4\theta E \frac{W^2\bar{W}^2}{\mathcal{D}^2 W^2 \bar{\mathcal{D}}^2 \bar{W}^2} \mathcal{D}^\alpha W_\alpha .\end{aligned}$$

- ▶ And in components

$$\begin{aligned}e^{-1}\mathcal{L}_{NEW} = & -\frac{1}{4}F_{mn}F^{mn} - i\bar{\lambda}\bar{\sigma}^m\mathcal{D}_m\lambda + \frac{1}{2}\mathbf{D}^2 \\ & -\xi\mathbf{D} + \xi\mathcal{O}(\lambda, \bar{\lambda}) .\end{aligned}$$

- ▶ The gravitino is not charged under the $U(1)$.

The new term coupled to pure supergravity

- ▶ In superspace the coupling is

$$\mathcal{L} = -3 \left(\int d^2\Theta 2\mathcal{E} \mathcal{R} + c.c. \right) + \left(\int d^2\Theta 2\mathcal{E} W_0 + c.c. \right) + \mathcal{L}_{NEW}.$$

- ▶ And in components (on-shell)

$$\begin{aligned} e^{-1} \mathcal{L} \Big|_{\lambda=0} = & -\frac{1}{2} R + \frac{1}{2} \epsilon^{klmn} (\bar{\psi}_k \bar{\sigma}_l \mathcal{D}_m \psi_n - \psi_k \sigma_l \mathcal{D}_m \bar{\psi}_n) \\ & - \frac{1}{4} F_{mn} F^{mn} - \left(\frac{1}{2} \xi^2 - 3 |W_0|^2 \right) \\ & - \bar{W}_0 \psi_a \sigma^{ab} \psi_b - W_0 \bar{\psi}_a \bar{\sigma}^{ab} \bar{\psi}_b. \end{aligned}$$

- ▶ A gravitino mass is allowed and R-symmetry is not gauged!

A KKLT-type uplift!

→ For a generic K and W the scalar potential becomes

$$\mathcal{V} = \mathcal{V}_{\text{Standard}}^{\text{SUGRA}} + \frac{\xi^2}{2} e^{2K/3}.$$

→ This uplift usually constructed in SUGRA with non-linear local SUSY realization. *Ferrara, Kallosh, Linde '14, Kallosh, Wrase '14, Bandos, Heller, Kuzenko, Martucci, Sorokin '16*

→ Here this construction is straightforward and achieved without introducing non-linear realizations off-shell.

Cribiori, Tournoy, FF, Van Proeyen '17

Thank you