

Variational characterization of circles in low dimensions

Josef Šilhan (joint with M. Dafinger)

Masaryk University, Brno, Czech Republic

43th Winter School “Geometry and Physics”, Srní, 14.1.2023

Plan

- Variational problem for circles in \mathbb{R}^n
- Inverse problem and its solution in dim 3

Variational problem for circles in \mathbb{R}^n

Inverse problem and its solution in dim 3

— Let $X(t) : I \rightarrow \mathbb{R}^n$ be a parametrized curve on some interval I . Its velocity vector, acceleration vector etc. are denoted by

$$U = X' = \frac{d}{dt}X, \quad U' = X'' = \frac{d^2}{dt^2}X, \quad U'' = \frac{d^3}{dt^3}X$$

where $(\)' = \frac{d}{dt} = U^r \nabla_r (\)$. Here $\nabla_a = \frac{\partial}{\partial x^a}$ is the Levi-Civita connection of the Euclidean metric.

— Variational approach: given a Lagrangian

$$\mathcal{L} = \mathcal{L}(t, X, U, U', \dots, U^{(k)}) : I \rightarrow \mathbb{R},$$

find extremal curves – or just critical curves – of the integral

$$\int_{t_1}^{t_2} \mathcal{L} dt, \quad I = [t_1, t_2]$$

among all curves with fixed endpoints $X(t_i)$ and derived vectors $U^{(j)}(t_i)$ at endpoints, $0 \leq j \leq k$.

— If the integral $\int_{t_1}^{t_2} \mathcal{L} dt$ is independent on reparametrization, then any reparametrization of a critical curve is again critical.

— Let $X(t) : I \rightarrow \mathbb{R}^n$ be a parametrized curve on some interval I . Its velocity vector, acceleration vector etc. are denoted by

$$U = X' = \frac{d}{dt}X, \quad U' = X'' = \frac{d^2}{dt^2}X, \quad U'' = \frac{d^3}{dt^3}X$$

where $(\)' = \frac{d}{dt} = U^r \nabla_r (\)$. Here $\nabla_a = \frac{\partial}{\partial x^a}$ is the Levi-Civita connection of the Euclidean metric.

— Variational approach: given a Lagrangian

$$\mathcal{L} = \mathcal{L}(t, X, U, U', \dots, U^{(k)}) : I \rightarrow \mathbb{R},$$

find extremal curves – or just critical curves – of the integral

$$\int_{t_1}^{t_2} \mathcal{L} dt, \quad I = [t_1, t_2]$$

among all curves with fixed endpoints $X(t_j)$ and derived vectors $U^{(j)}(t_j)$ at endpoints, $0 \leq j \leq k$.

— If the integral $\int_{t_1}^{t_2} \mathcal{L} dt$ is independent on reparametrization, then any reparametrization of a critical curve is again critical.

— Let $X(t) : I \rightarrow \mathbb{R}^n$ be a parametrized curve on some interval I . Its velocity vector, acceleration vector etc. are denoted by

$$U = X' = \frac{d}{dt}X, \quad U' = X'' = \frac{d^2}{dt^2}X, \quad U'' = \frac{d^3}{dt^3}X$$

where $(\)' = \frac{d}{dt} = U^r \nabla_r (\)$. Here $\nabla_a = \frac{\partial}{\partial x^a}$ is the Levi-Civita connection of the Euclidean metric.

— Variational approach: given a Lagrangian

$$\mathcal{L} = \mathcal{L}(t, X, U, U', \dots, U^{(k)}) : I \rightarrow \mathbb{R},$$

find extremal curves – or just critical curves – of the integral

$$\int_{t_1}^{t_2} \mathcal{L} dt, \quad I = [t_1, t_2]$$

among all curves with fixed endpoints $X(t_j)$ and derived vectors $U^{(j)}(t_j)$ at endpoints, $0 \leq j \leq k$.

— If the integral $\int_{t_1}^{t_2} \mathcal{L} dt$ is independent on reparametrization, then any reparametrization of a critical curve is again critical.

— Which families of curves are *variational* in the sense that they form the family of critical curves of suitable \mathcal{L} ? And what is the lowest order for such \mathcal{L} ?

- ▶ This is generally a difficult question
- ▶ Classical example – straight lines are variational:
 - ▶ $\mathcal{L} = \langle U, U \rangle \rightsquigarrow$ energy functional (not parametrization invariant)
 - ▶ $\mathcal{L} = \langle U, U \rangle^{1/2} \rightsquigarrow$ length functional (parametrization invariant)

— Our problem: is the family of all circles – or of all conformal circles – variational? If so, what is the lowest order Lagrangian?

- ▶ find the required Lagrangian (optimally parametrization invariant)
- ▶ or show its nonexistence

— Which families of curves are *variational* in the sense that they form the family of critical curves of suitable \mathcal{L} ? And what is the lowest order for such \mathcal{L} ?

- ▶ This is generally a difficult question
- ▶ Classical example – straight lines are variational:
 - ▶ $\mathcal{L} = \langle U, U \rangle \rightsquigarrow$ energy functional (not parametrization invariant)
 - ▶ $\mathcal{L} = \langle U, U \rangle^{1/2} \rightsquigarrow$ length functional (parametrization invariant)

— Our problem: is the family of all circles – or of all conformal circles – variational? If so, what is the lowest order Lagrangian?

- ▶ find the required Lagrangian (optimally parametrization invariant)
- ▶ or show its nonexistence

— The usual torsion τ of curves (more or less) solves the problem on \mathbb{R}^3 with parametrization independent Lagrangian,

$$\mathcal{L} = \tau(U, U', U'') \langle U, U' \rangle^{1/2} = \frac{\epsilon(U, U', U'')}{G(U, U')} \langle U, U' \rangle^{1/2}$$

where $G(U, U')$ is the Gram matrix of U and U' :

- ▶ critical curves are exactly all circles in \mathbb{R}^3 (but not straight lines)
- ▶ a geometrical explanation is unclear
- ▶ N. Thamwattana, J.A. McCoy, J.M. Hill. *Energy density functions for protein structures*, Q. J. Mech. Appl. Math. 61(3):431451 (2008).

— But is this the lowest order Lagrangian for circles in \mathbb{R}^3 ?

- ▶ No.

— Circles are characterized by constant curvature κ and zero torsion τ . Including also straight lines, we have two ODEs

$$\kappa = \left(\frac{G(U, U')}{\langle U, U \rangle^3} \right)^{\frac{1}{2}} \in \mathbb{R} \quad \text{and} \quad U'' \in \text{span}\{U, U'\}.$$

Put $u = \sqrt{\langle U, U \rangle}$. Expanding the latter condition, we obtain the system of n ODEs

$$\underbrace{U'' + u^{-2}[-\langle U, U'' \rangle + 3u^{-2}\langle U, U' \rangle^2]U - 3u^{-2}\langle U, U' \rangle U'}_{E(U, U', U'')} = 0.$$

— There are many other systems of ODEs $\bar{E}_a = 0$ characterizing (conformal) parameter-independent circles. BUT the system $E_a = 0$ is in a suitable sense *nondegenerate* $\leadsto \bar{E}_a = 0$ has the form

$$\bar{E}_a = M^r{}_a E_r \quad \text{for a multiplier} \quad M^b{}_a = M^b{}_a(t, X, U, U', U'').$$

— Circles are characterized by constant curvature κ and zero torsion τ . Including also straight lines, we have two ODEs

$$\kappa = \left(\frac{G(U, U')}{\langle U, U \rangle^3} \right)^{\frac{1}{2}} \in \mathbb{R} \quad \text{and} \quad U'' \in \text{span}\{U, U'\}.$$

Put $u = \sqrt{\langle U, U \rangle}$. Expanding the latter condition, we obtain the system of n ODEs

$$\underbrace{U'' + u^{-2}[-\langle U, U'' \rangle + 3u^{-2}\langle U, U' \rangle^2]U - 3u^{-2}\langle U, U' \rangle U'}_{E(U, U', U'')} = 0.$$

— There are many other systems of ODEs $\bar{E}_a = 0$ characterizing **(conformal) parameter-independent circles**. BUT the system $E_a = 0$ is in a suitable sense *nondegenerate* $\rightsquigarrow \bar{E}_a = 0$ has the form

$$\bar{E}_a = M^r{}_a E_r \quad \text{for a multiplier } M^b{}_a = M^b{}_a(t, X, U, U', U'').$$

Variational problem for circles in \mathbb{R}^n

Inverse problem and its solution in dim 3

— Variationality of the system $\bar{E}_a = 0$ means it is the Euler-Lagrange equation of a Lagrangian

$$\mathcal{L} = \mathcal{L}(t, X, U, U', U''),$$

i.e. it has the form

$$\bar{E}_a = \frac{\partial \mathcal{L}}{\partial X_a} - \frac{d}{dt} \frac{\partial \mathcal{L}}{\partial U_a} + \frac{d^2}{dt^2} \frac{\partial \mathcal{L}}{\partial U'_a} - \frac{d^3}{dt^3} \frac{\partial \mathcal{L}}{\partial U''_a} = 0.$$

— Equations of this form are characterized by Helmholtz conditions

$$\frac{\partial \bar{E}_a}{\partial X^b} - \frac{\partial \bar{E}_b}{\partial X^a} - \frac{1}{2} \frac{d}{dt} \left(\frac{\partial \bar{E}_a}{\partial U^b} - \frac{\partial \bar{E}_b}{\partial U^a} \right) + \frac{1}{4} \frac{d^3}{dt^3} \left(\frac{\partial \bar{E}_a}{\partial U''^b} - \frac{\partial \bar{E}_b}{\partial U''^a} \right) = 0,$$

$$\frac{\partial \bar{E}_a}{\partial U^b} + \frac{\partial \bar{E}_b}{\partial U^a} - \frac{d}{dt} \left(\frac{\partial \bar{E}_a}{\partial U'^b} + \frac{\partial \bar{E}_b}{\partial U'^a} \right) = 0,$$

$$\frac{\partial \bar{E}_a}{\partial U'^b} - \frac{\partial \bar{E}_b}{\partial U'^a} - \frac{3}{2} \frac{d}{dt} \left(\frac{\partial \bar{E}_a}{\partial U''^b} - \frac{\partial \bar{E}_b}{\partial U''^a} \right) = 0,$$

$$\frac{\partial \bar{E}_a}{\partial U''^b} + \frac{\partial \bar{E}_b}{\partial U''^a} = 0.$$

— Variationality of the system $\bar{E}_a = 0$ means it is the Euler-Lagrange equation of a Lagrangian

$$\mathcal{L} = \mathcal{L}(t, X, U, U', U''),$$

i.e. it has the form

$$\bar{E}_a = \frac{\partial \mathcal{L}}{\partial X_a} - \frac{d}{dt} \frac{\partial \mathcal{L}}{\partial U_a} + \frac{d^2}{dt^2} \frac{\partial \mathcal{L}}{\partial U'_a} - \frac{d^3}{dt^3} \frac{\partial \mathcal{L}}{\partial U''_a} = 0.$$

— Equations of this form are characterized by Helmholtz conditions

$$\frac{\partial \bar{E}_a}{\partial X^b} - \frac{\partial \bar{E}_b}{\partial X^a} - \frac{1}{2} \frac{d}{dt} \left(\frac{\partial \bar{E}_a}{\partial U^b} - \frac{\partial \bar{E}_b}{\partial U^a} \right) + \frac{1}{4} \frac{d^3}{dt^3} \left(\frac{\partial \bar{E}_a}{\partial U''^b} - \frac{\partial \bar{E}_b}{\partial U''^a} \right) = 0,$$

$$\frac{\partial \bar{E}_a}{\partial U^b} + \frac{\partial \bar{E}_b}{\partial U^a} - \frac{d}{dt} \left(\frac{\partial \bar{E}_a}{\partial U'^b} + \frac{\partial \bar{E}_b}{\partial U'^a} \right) = 0,$$

$$\frac{\partial \bar{E}_a}{\partial U'^b} - \frac{\partial \bar{E}_b}{\partial U'^a} - \frac{3}{2} \frac{d}{dt} \left(\frac{\partial \bar{E}_a}{\partial U''^b} - \frac{\partial \bar{E}_b}{\partial U''^a} \right) = 0,$$

$$\frac{\partial \bar{E}_a}{\partial U''^b} + \frac{\partial \bar{E}_b}{\partial U''^a} = 0.$$

— We need to find (or to prove nonexistence) of a multiplier M^b_a such that $\bar{E}_a = M^r_a E_r$ satisfies Helmholtz conditions. This is generally extremely difficult.

— Our setup: assume $M^b_a = M^b_a(t, X, U, U')$. Then the simplest Helmholtz condition simplifies the problem significantly in low dimensions:

- ▶ Dimension $n = 2$: such multiplier does not exist \rightsquigarrow (conformal) circles in the plane are not variational as unparametrized curves.
- ▶ Dimension $n = 3$: the multiplier has the form

$$M^b_a = \varphi U^r \epsilon_r^b{}_a, \quad \varphi = \varphi(t, X, U, U')$$

where ϵ_{abc} is the Euclidean volume form on \mathbb{R}^3 . It remains to determine the function φ .

- ▶ Dimension $n \geq 4$: unclear.

— Remaining Helmholtz cond's for $n = 3 \rightsquigarrow \varphi$ is a constant.

- We need to find (or to prove nonexistence) of a multiplier M^b_a such that $\bar{E}_a = M^r_a E_r$ satisfies Helmholtz conditions. This is generally extremely difficult.
- Our setup: assume $M^b_a = M^b_a(t, X, U, U')$. Then the simplest Helmholtz condition simplifies the problem significantly in low dimensions:
 - ▶ Dimension $n = 2$: such multiplier does not exist \rightsquigarrow (conformal) circles in the plane are not variational as unparametrized curves.
 - ▶ Dimension $n = 3$: the multiplier has the form

$$M^b_a = \varphi U^r \epsilon_r^b{}_a, \quad \varphi = \varphi(t, X, U, U')$$

where ϵ_{abc} is the Euclidean volume form on \mathbb{R}^3 . It remains to determine the function φ .

- ▶ Dimension $n \geq 4$: unclear.

— Remaining Helmholtz cond's for $n = 3 \rightsquigarrow \varphi$ is a constant.

- We need to find (or to prove nonexistence) of a multiplier M^b_a such that $\bar{E}_a = M^r_a E_r$ satisfies Helmholtz conditions. This is generally extremely difficult.
- Our setup: assume $M^b_a = M^b_a(t, X, U, U')$. Then the simplest Helmholtz condition simplifies the problem significantly in low dimensions:
 - ▶ Dimension $n = 2$: such multiplier does not exist \rightsquigarrow (conformal) circles in the plane are not variational as unparametrized curves.
 - ▶ Dimension $n = 3$: the multiplier has the form

$$M^b_a = \varphi U^r \epsilon_r^b{}_a, \quad \varphi = \varphi(t, X, U, U')$$

where ϵ_{abc} is the Euclidean volume form on \mathbb{R}^3 . It remains to determine the function φ .

- ▶ Dimension $n \geq 4$: unclear.
- Remaining Helmholtz cond's for $n = 3 \rightsquigarrow \varphi$ is a constant.

— We found the variational system of ODEs

$$\bar{E}_a = \epsilon(U, U'', -)_a - 3u^2 \langle U, U' \rangle \epsilon(U, U', -)_a = 0$$

for (conformal) unparametrized circles in \mathbb{R}^3 .

— What is the corresponding Lagrangian $\mathcal{L} = \mathcal{L}(t, X, U, U')$? The answer starts with the Vainberg-Tonti formula

$$\text{Lagrangian} := \int_0^1 \bar{E}(sU + V, sU', sU'')_a X^a ds$$

for any $V \in \mathbb{R}^3$.

- ▶ the case $V = 0$ is the usual Vainberg-Tonti formula but this integral does not exist (it is infinite).
- ▶ we can modify the Lagrangian by total derivative term
- ▶ then the integration is manageable by hand
- ▶ after some other technicalities we obtain the result.

— We found the variational system of ODEs

$$\bar{E}_a = \epsilon(U, U'', -)_a - 3u^2 \langle U, U' \rangle \epsilon(U, U', -)_a = 0$$

for (conformal) unparametrized circles in \mathbb{R}^3 .

— What is the corresponding Lagrangian $\mathcal{L} = \mathcal{L}(t, X, U, U')$? The answer starts with the Vainberg-Tonti formula

$$\text{Lagrangian} := \int_0^1 \bar{E}(sU + V, sU', sU'')_a X^a ds$$

for any $V \in \mathbb{R}^3$.

- ▶ the case $V = 0$ is the usual Vainberg-Tonti formula but this integral does not exist (it is infinite).
- ▶ we can modify the Lagrangian by total derivative term
- ▶ then the integration is manageable by hand
- ▶ after some other technicalities we obtain the result.

— Assume $\|V\| = 1$. After integration and further manipulation, we obtain

$$\mathcal{L} := \frac{\epsilon(V, U, U')}{\|U\|(\|U\| + \langle U, V \rangle)}.$$

One can directly verify that solutions of the EL-equations are exactly conformal circles.

— We found two Lagrangians for (conformal) unparametrized circles in \mathbb{R}^3 but both have singularities:

- ▶ torsion τ excludes straight lines
- ▶ the Lagrangian \mathcal{L} excludes tangent directions $-V$

— Final comments/questions:

- ▶ is there a globally defined Lagrangian for conformal circles?
- ▶ do we have something similar for non-flat metrics?
- ▶ what about higher dimensions?

— Assume $\|V\| = 1$. After integration and further manipulation, we obtain

$$\mathcal{L} := \frac{\epsilon(V, U, U')}{\|U\|(\|U\| + \langle U, V \rangle)}.$$

One can directly verify that solutions of the EL-equations are exactly conformal circles.

— We found two Lagrangians for (conformal) unparametrized circles in \mathbb{R}^3 but both have singularities:

- ▶ torsion τ excludes straight lines
- ▶ the Lagrangian \mathcal{L} excludes tangent directions $-V$

— Final comments/questions:

- ▶ is there a globally defined Lagrangian for conformal circles?
- ▶ do we have something similar for non-flat metrics?
- ▶ what about higher dimensions?

— Assume $\|V\| = 1$. After integration and further manipulation, we obtain

$$\mathcal{L} := \frac{\epsilon(V, U, U')}{\|U\|(\|U\| + \langle U, V \rangle)}.$$

One can directly verify that solutions of the EL-equations are exactly conformal circles.

— We found two Lagrangians for (conformal) unparametrized circles in \mathbb{R}^3 but both have singularities:

- ▶ torsion τ excludes straight lines
- ▶ the Lagrangian \mathcal{L} excludes tangent directions $-V$

— Final comments/questions:

- ▶ is there a globally defined Lagrangian for conformal circles?
- ▶ do we have something similar for non-flat metrics?
- ▶ what about higher dimensions?

— Assume $\|V\| = 1$. After integration and further manipulation, we obtain

$$\mathcal{L} := \frac{\epsilon(V, U, U')}{\|U\|(\|U\| + \langle U, V \rangle)}.$$

One can directly verify that solutions of the EL-equations are exactly conformal circles.

— We found two Lagrangians for (conformal) unparametrized circles in \mathbb{R}^3 but both have singularities:

- ▶ torsion τ excludes straight lines
- ▶ the Lagrangian \mathcal{L} excludes tangent directions $-V$

— Final comments/questions:

- ▶ is there a globally defined Lagrangian for conformal circles?
- ▶ do we have something similar for non-flat metrics?
- ▶ what about higher dimensions?

Thank you for your attention!