## Non-Abelian Gerbes with Connections and with an explicit cocycle description

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Gauge fields are everywhere:

- Gauge fields in the Standard Model and in Yang-Mills theory
- Chern–Simons (matter) theories
  - Importantly, M2-brane models: BLG- and ABJM-model
- Kalb-Ramond B-field in string theory
- B-field and connection in heterotic supergravity
- Tensor hierarchies in gauged supergravity
- T-duality: *B*-fields on top of circle bundles with connection
- Interacting M5-branes:
  - self-dual strings
  - (2,0)-theory

# Motivation: M-theory





"One (superst)ring to rule them all ..."

# Motivation: (2, 0)-theory



The picture so far:

- M5 branes interact via M2 branes, boundaries = self-dual strings
- 6D SCFT from M-theory with (2,0) SUSY
- No Lagrangian exists irreducibly quantum?
- Field content: (2,0) tensor multiplet includes 2-form *B*-field
- This B-field satisfies the self-duality equation:
   H := dB = \*H

#### What we need:

- to go beyond connections on principal bundles
- a framework to describe higher connection forms, e.g. for parallel transport of extended objects

**Good news!** There is such a framework – higher gauge theory.

#### Categories, Higher Structures, and All That

# What are categories?

A category  $\mathscr{C} = (\mathscr{C}_0 \models \mathscr{C}_1)$  consists of

- a collection  $\mathscr{C}_0$  of objects,
- a collection  $\mathscr{C}_1$  of morphisms between objects such that



When all morphisms are invertible  $\implies$  groupoid. Group = one-object groupoid, BG = (G  $\implies$  \*). Morphisms between categories = functors.

# Fun with functors

Principal bundles = functors from  $\check{C}(U)$  to BG.

• Cover  $\mathcal{U} = \sqcup_a U_a$  of a manifold M yields Čech groupoid  $\check{\mathcal{C}}(\mathcal{U})$ :

$$(x,a) \xleftarrow{(x,a,b)}{(x,b)} (x,b) \xleftarrow{(x,b,c)}{(x,c)} (x,c)$$



Transition functions  $g_{ab}$ , cocycle cond.  $g_{ab}g_{bc} = g_{ac}$ cobndries.:  $g_{ab}\gamma_b = \gamma_a \tilde{g}_{ab}$ 

# Categories not enough!

Parallel transport of extended objects: replace paths by (higher-dimensional) surfaces.

Consider the parallel transport of one-dimensional strings along a two-dimensional surface



This means that

- $(g_2g_1)(g'_2g'_1) = (g_2g'_2)(g_1g'_1).$
- Setting  $g_2 = 1$  and  $g'_1 = 1$  yields  $g_1g'_2 = g'_2g_1$ .
- $\Rightarrow$  Forces G to be abelian = Eckmann-Hilton argument.

#### Problem!

#### String/M- theory requires the framework of higher category theory.

### 2-categories to the rescue

Recall a category  $\mathscr{C} = (\mathscr{C}_1 \rightrightarrows \mathscr{C}_0)$  has  $\mathscr{C}_1$  as the collection of morphisms. To construct a (strict) 2-category

- take *C*<sub>1</sub> to be a category itself, i.e. introduce morphisms between morphisms = 2-morphisms,
- introduce associative horizontal composition ⊗ between 2-morphisms (the other being vertical ○) such that the following interchange law holds

$$(\eta_2\circ\eta_1)\otimes(\eta_2'\circ\eta_1')=(\eta_2\otimes\eta_2')\circ(\eta_1\otimes\eta_1')$$



Morphisms of 2-categories = 2-functors.

A mathematical structure ("Bourbaki-style") consists of

• Sets • Structure Functions • Structure Equations Categorification: "adding morphisms between morphisms"

 $\mathsf{Sets} \to \mathsf{Categories}$ 

Structure Functions  $\rightarrow$  Structure Functors Structure Equations  $\rightarrow$  Structure Isomorphisms

Example: Group(oid)  $\rightarrow$  2-Group(oid)

Note: Process not unique, variants: weak/strict/...

# Strict 2-groups

Only need strict Lie 2-groups = crossed modules of Lie groups:

- a pair of Lie groups (G, H) with an automorphism action  $\vartriangleright$  of G on H
- $\bullet$  a morphism of Lie groups  $t: H \rightarrow G$  such that

$$t(g 
ho h_1) = gt(h_1)g^{-1}$$
 and  $t(h_1) 
ho h_2 = h_1h_2h_1^{-1}$ 

for all  $g \in G$  and for all  $h_1, h_2 \in H$ .

- Written as  $\mathcal{G} := (\mathsf{H} \stackrel{\mathsf{t}}{\longrightarrow} \mathsf{G}, \rhd).$
- Play the role of structure groups for principal 2-bundles a.k.a. gerbes.

#### Examples:

- $(U(1) \longrightarrow 1) \cong BU(1)$  underlies an abelian gerbe.
- $(1 \hookrightarrow G, id)$  underlies a principal G-bundle (flat?)
- $\mathcal{L}G := (L_0G \hookrightarrow P_0G, Ad)$

Equivalences between crossed modules given by flippable butterflies:



where E is a Lie group,  $\lambda_i, \gamma_i$  are morphisms of Lie groups and both diagonals are short exact sequences.

# **Butterflies**

Recall crossed modules  $\mathcal{L}G = (L_0G \hookrightarrow P_0G, Ad)$  and  $(\mathbb{1} \hookrightarrow G, id)$ . Are they equivalent?





Here,  $\flat : P_0 G \rightarrow G$  is the endpoint evaluation map. They are **equivalent**!

# Gerbes without connections

Principal  $\mathcal{G}$ -bundle = 2-functor from  $\check{\mathcal{C}}(\mathcal{U})$  to  $B\mathcal{G}$ .

• 
$$\mathcal{G} := (\mathsf{H} \stackrel{\mathsf{t}}{\longrightarrow} \mathsf{G}, \rhd)$$

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• Trivially regard  $\check{\mathcal{C}}(\mathcal{U})$  as a 2-groupoid by adding identity 2-morphisms

• Transition functions  $(g_{ab}, h_{abc})$  satisfying cocycle conditions:

$$t(h_{abc})g_{ab}g_{bc} = g_{ac} \;, \quad h_{acd}h_{abc} = h_{abd}(g_{ab} \rhd h_{bcd})$$

• Equivalent when related by coboundaries  $(g_a, h_{ab})$ :

$$g_ag_{ab}'={
m t}(h_{ab})g_{ab}g_b\;,\quad h_{ac}h_{abc}=(g_a\rhd\,h_{abc}')h_{ab}(g_{ab}\rhd\,h_{bc})$$

The Local Picture:  $L_{\infty}$ -algebras and Higher Connections

# Crossed modules of Lie algebras

Differentiating a crossed module of Lie groups yields a crossed module of Lie algebras:

- given by a pair  $(\mathfrak{g}, \mathfrak{h})$  of Lie algebras,
- automorphism action  $\triangleright$  of  $\mathfrak{g}$  on  $\mathfrak{h}$ ,
- $\bullet\,$  morphism of Lie algebras  $t:\mathfrak{h}\to\mathfrak{g}$  such that

 $\mathsf{t}(X \vartriangleright Y_1) \ := [X, \mathsf{t}(Y_1)] \quad \text{and} \quad \mathsf{t}(Y_1) \vartriangleright Y_2 \ := [Y_1, Y_2] \ .$ 

- Can be viewed as a (strict 2-term)  $L_{\infty}$ -algebra:
  - $\mathbb{Z}$ -graded vector space  $\mathfrak{L} = \underbrace{\mathfrak{L}_{-1}}_{h} \oplus \underbrace{\mathfrak{L}_{0}}_{\mathfrak{g}}$ ,
  - with differential  $\mathsf{t}=\mu_1:\mathfrak{L}_{-1}\to\mathfrak{L}_{\mathsf{0}},$
  - with graded anti-symmetric Lie bracket μ<sub>2</sub> : ℒ<sub>i</sub> × ℒ<sub>j</sub> → ℒ<sub>i+j</sub>, (α<sub>1</sub>, α<sub>2</sub>) ↦ [α<sub>1</sub>, α<sub>2</sub>], (α<sub>1</sub>, β<sub>1</sub>) ↦ α<sub>1</sub> ▷ β<sub>1</sub>,
  - such that  $\mu_1$  is a derivation and  $\mu_2$  satisfies a graded Jacobi identity + higher maps

- Generalizations of (differential graded) Lie algebras
- Come with their own gauge theory:
  - elements  $a \in \mathfrak{L}_1$  of degree 1 are generalized gauge potentials
  - curvature:  $f := \mu_1(a) + \frac{1}{2}\mu_2(a,a) + \frac{1}{3!}\mu_3(a,a,a) + \dots \in \mathfrak{L}_2$ 
    - Homotopy Maurer-Cartan equation: f = 0
  - Bianchi identity
    - $\mu_1(f) \mu_2(f, a) + \frac{1}{2}\mu_3(f, a, a) \frac{1}{3!}\mu_4(f, a, a, a) + \cdots = 0$
  - $\bullet~$  Elements  $\in \mathfrak{L}_0$  parametrize generalized gauge transformations.
- Ordinary gauge theory on manifold *M* modelled on Ω<sup>●</sup>(*M*) ⊗ g
- Higher gauge theory modelled on  $\Omega^{\bullet}(M) \otimes \mathfrak{L}$ .

# $L_{\infty}$ -algebras

The tensor product  $\hat{\mathfrak{L}} = \Omega^{ullet}(M) \otimes \mathfrak{L}$  is an  $L_{\infty}$ -algebra as well

- connection forms = gauge potentials  $A + B + \cdots \in \hat{\mathfrak{L}}_1 = (\Omega^1(M) \otimes \mathfrak{L}_0) \oplus (\Omega^2(M) \otimes \mathfrak{L}_{-1}) \oplus \cdots$
- higher products  $\hat{\mu}_1 = d + \mu_1$ ,  $\mu_2$ ,  $\mu_3$ , ...
- curvatures

$$F = dA + \frac{1}{2}\mu_2(A, A) + \mu_1(B) \rightarrow \text{ fake curvature}$$
$$H = dB + \mu_2(A, B) - \frac{1}{3!}\mu_3(A, A, A) + \dots$$

 Infinitesimal gauge transformations can be obtained from the definition of curvatures as partially flat homotopies. They are parametrised by
 α + Λ ∈ L̂<sub>0</sub> = (Ω<sup>0</sup>(M) ⊗ L̂<sub>0</sub>) ⊕ (Ω<sup>1</sup>(M) ⊗ L̂<sub>-1</sub>) ⊕ · · · .
 **Problem:** consistency requires fake flatness  $F = dA + \frac{1}{2}\mu_2(A, A) + \mu_1(B) = 0$ 

- For  $\mu_3 \neq 0$  infinitesimal gauge transformations do not close:  $[\delta_{c_0}, \delta_{c_1}] = \delta_{[c_0, c_1]} + \mu_3(F, c_0, c_1).$
- For  $\mu_3 = 0$  finite gauge transformations do not compose.
- Self-duality  $H = \star H$  requires F = 0 for covariance.
- Parallel transport requires F = 0 for reparametrisation invariance.
- Can gauge away any non-Abelian part  $\implies$  theory Abelian.

#### Problem!

**Solution:** adjust the definition of curvatures:  $H \rightarrow H + \kappa(A, F)$  where  $\kappa$  is called the adjustment datum. Examples:

• 
$$\mathcal{L}\mathfrak{g} = (L_0\mathfrak{g} \hookrightarrow P_0\mathfrak{g})$$
:

- connection given by 1- and 2-forms  $(A, B) \in \Omega^1(M, P_0\mathfrak{g}) \oplus \Omega^2(M, L_0\mathfrak{g}).$
- curvature forms

$$\begin{split} F &:= \mathrm{d}A + \tfrac{1}{2}[A,A] + \mathrm{t}(B) \in \Omega^2(M,P_0\mathfrak{g}) \ , \\ H &:= \mathrm{d}B + A \rhd B - \kappa(A,F) \in \Omega^3(M,L_0\mathfrak{g}) \ , \end{split}$$

where  $\kappa(A, F) = (1 - \wp \cdot \flat)([A, F])$  for an arbitrary fixed path  $\wp \in \mathscr{C}^{\infty}([0, 1], \mathbb{R})$  such that  $\wp(0) = 0$  and  $\wp(1) = 1$ . • String Lie 2-algebra:  $\mathfrak{string}(G) = (L_0\mathfrak{g} \oplus \mathfrak{u}(1) \to P_0\mathfrak{g})$ •  $\hat{\kappa}(A, F) = \left(\kappa(A, F), \frac{i}{2\pi} \int_0^1 dr \langle A_r, F \rangle\right)$ 

#### The Global Picture: Non-Abelian Gerbes with Connections

# Non-Abelian gerbes with connections

Putting everything together to construct principal 1- and 2bundles with connections.

• Manifold M with cover  $\mathcal{U} = (U_a)$ 

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• Lie 2-group and 2-algebra: crossed modules  $\mathsf{H}\to\mathsf{G}$  and  $\mathfrak{h}\to\mathfrak{g}$  with adjustment datum  $\kappa$ 

Object	Principal G-bundle	Principal (H $\xrightarrow{t}$ G)-bundle
Cochains	(g <sub>ab</sub> ) valued in G	$(g_{ab})$ valued in G, $(h_{abc})$ valued in H
Cocycle	$g_{ab}g_{bc}=g_{ac}$	$t(h_{abc})g_{ab}g_{bc}=g_{ac} \ h_{acd}h_{abc}=h_{abd}(g_{ab}arprophoh_{bcd})$
Coboundary	$g_a g_{ab}' = g_{ab} g_b$	$egin{array}{llllllllllllllllllllllllllllllllllll$
gauge pot.	${oldsymbol{\mathcal{A}}}_{{oldsymbol{a}}}\in \Omega^1(U_{{oldsymbol{a}}})\otimes \mathfrak{g}$	${\mathcal A}_{{\mathfrak a}}\in \Omega^1(U_{{\mathfrak a}})\otimes {\mathfrak g},\ {\mathcal B}_{{\mathfrak a}}\in \Omega^2(U_{{\mathfrak a}})\otimes {\mathfrak h}$
Curvature	$F_a = \mathrm{d}A_a + \frac{1}{2}[A_a, A_a]$	$ \begin{aligned} & F_{\boldsymbol{\partial}} = \mathrm{d}\mathcal{A}_{\boldsymbol{\partial}} + \frac{1}{2}[\mathcal{A}_{\boldsymbol{\partial}}, \mathcal{A}_{\boldsymbol{\partial}}] + \mathrm{t}(\mathcal{B}) \\ & H_{\boldsymbol{\partial}} = \mathrm{d}\mathcal{B}_{\boldsymbol{\partial}} + \mathcal{A}_{\boldsymbol{\partial}} \rhd \mathcal{B}_{\boldsymbol{\partial}} - \kappa(\mathcal{A}_{\boldsymbol{\partial}}, \mathcal{F}_{\boldsymbol{\partial}}) \end{aligned} $
Gauge trafos	$ ilde{A}_{a}:=g_{a}^{-1}A_{a}g_{a}+g_{a}^{-1}\mathrm{d}g_{a}$	$\begin{split} \tilde{A}_{a} &:= g_{a}^{-1} A_{a} g_{a} + g_{a}^{-1} \mathrm{d} g_{a} - t(\Lambda_{a}) \\ \tilde{B}_{a} &:= g_{a}^{-1} \rhd B_{a} + \tilde{A}_{a} \rhd \Lambda_{a} + \mathrm{d} \Lambda_{a} + \frac{1}{2} [\Lambda_{a}, \Lambda_{a}] - \kappa(g_{a}, F_{a}) \end{split}$

Consider a surjective submersion  $\sigma: Y \to M$  and an adjusted crossed module of Lie groups  $\mathcal{G} := (H \stackrel{t}{\longrightarrow} G, \rhd, \kappa)$  with the corresponding crossed module of Lie algebras  $(\mathfrak{h} \stackrel{t}{\longrightarrow} \mathfrak{g}, \rhd, \kappa)$ . An adjusted cocycle for a principal  $\mathcal{G}$ -bundle over M is given by the data

$$\begin{array}{l} h \ \in \ \mathscr{C}^{\infty}(Y^{[3]},\mathsf{H}) \ , \\ (g,\Lambda) \ \in \ \mathscr{C}^{\infty}(Y^{[2]},\mathsf{G}) \oplus \Omega^{1}(Y^{[2]},\mathfrak{h}) \ , \\ (A,B) \ \in \ \Omega^{1}(Y^{[1]},\mathfrak{g}) \oplus \Omega^{2}(Y^{[1]},\mathfrak{h}) \ , \end{array}$$

such that

$$h_{ikl}h_{ijk} = h_{ijl}(g_{ij} \triangleright h_{jkl})$$
 and  $g_{ik} = t(h_{ijk})g_{ij}g_{jk}$   
hold for all appropriate  $(i, j, ...) \in Y^{[n]}$ .

In addition, we have the following cocycle conditions

$$\begin{split} \Lambda_{ik} &= \Lambda_{jk} + g_{jk}^{-1} \rhd \Lambda_{ij} - g_{ik}^{-1} \rhd \left(h_{ijk} \nabla_i h_{ijk}^{-1}\right) \,, \\ A_j &= g_{ij}^{-1} A_i g_{ij} + g_{ij}^{-1} \mathrm{d} g_{ij} - \mathsf{t}(\Lambda_{ij}) \,, \\ B_j &= g_{ij}^{-1} \rhd B_i + \mathrm{d} \Lambda_{ij} + A_j \rhd \Lambda_{ij} + \frac{1}{2} [\Lambda_{ij}, \Lambda_{ij}] - \kappa(g_{ij}, F_i) \end{split}$$

for all appropriate  $(i, j, ...) \in Y^{[n]}$ , where  $\kappa$  needs to satisfy the adjustment identity

$$\begin{aligned} (g_2^{-1}g_1^{-1}) &\rhd (h^{-1}(X \rhd h)) + g_2^{-1} \rhd \kappa(g_1, X) \\ &+ \kappa(g_2, g_1^{-1}Xg_1 - \mathsf{t}(\kappa(g_1, X))) - \kappa(\mathsf{t}(h)g_1g_2, X) = 0 \end{aligned}$$

for all  $g_1, g_2 \in G$ , for all  $h \in H$ , and for all  $X \in \mathfrak{g}$ .



- Higher structures not useless!
- Have now a proper definition of non-Abelian gerbes with connections and explicit examples.
- Can accommodate string structures!

To do:

- understand the origin of the adjustment datum
- construct physically relevant examples

#### Thank you for your attention!

#### "It pays off to have the right connections."

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