

# Algebraic structures and particles

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**String Field Theory**

**Spinning particles**

# String Field Theory

Second quantization of a string:

$$|\Psi\rangle \propto \int dk \Psi(k) |k\rangle$$

1. On-shell amplitudes: from punctured Riemann surfaces (expansion in Feynman diagrams).  
Moduli space:

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1. On-shell amplitudes: from punctured Riemann surfaces (expansion in Feynman diagrams).  
Moduli space:

- ▶ it is a cochain complex  $C_i \ni \nu_i$  with differential  $\partial$
- ▶ operation of joining two legs from different vertices and sewing two legs from the same vertex
- ▶ MC equation

$$\partial\nu_4 + \frac{1}{2}\{\nu_3, \nu_3\} + \underbrace{(\Delta\nu_6)}_{\text{loops}} = 0$$

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*Normally in a QFT, 1. and 2. are reversed.*

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*Knowledge of the gauge-fixed  $\mathcal{N} = 4$  **spinning** particle allows to calculate the 1-loop  $\Gamma[g_{\mu\nu}]$  of perturbative QG*

[Bastianelli et al.]

*Spinning particle* has a superworldline  $(\tau, \theta_j)$

Reparametrization invariance of superworldline  $\implies$  Supersymmetry algebra  $\mathfrak{g}$ :

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## $\mathcal{N} = 2$ BRST and quantization

- Needs a double cochain complex  $\mathcal{C}^{p,q} := \Lambda^p \mathfrak{g}^* \otimes \Lambda^q \mathfrak{g} \otimes C^\infty(\mathcal{M})$

inner derivations of  $\mathfrak{g}$  are  $\beta, \bar{\beta}, b$

outer derivations of  $\mathfrak{g}$  are  $\bar{\gamma}, \gamma, c$  with  
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- Pick vacuum  $\bar{\gamma} |0\rangle = \bar{\beta} |0\rangle = b |0\rangle = 0$ . States:

$$\underbrace{\omega_{k,m}(x, \psi)}_{\in \Gamma(E), E \xrightarrow{\pi} M \cong \mathcal{M}} \gamma^k \wedge \beta^m \wedge c |0\rangle \in \hat{\mathcal{C}}^{p,q}$$

In  $\hat{\mathcal{C}}^{p,q}$  one finds

$$A_\mu(x)\psi^\mu |0\rangle, C(x)\beta |0\rangle, \varphi(x)c_\beta |0\rangle$$

$$A_\mu^*(x)\psi^\mu c |0\rangle, C^*(x)c_\gamma |0\rangle, \varphi^*(x)\gamma |0\rangle$$

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The BRST nilpotent operator  $Q = e^i G_i + d_{CE}$  is:

$$Q = \frac{1}{2}cP^2 + \gamma\bar{q} + \bar{\gamma}q + \gamma\bar{\gamma}b \quad (1)$$

$Q$  acts on  $\mathcal{C}^{p,q}$  (and  $\hat{\mathcal{C}}^{p,q}$ ) by increasing  $p - q$  by  $+1$ .

# $\mathcal{N} = 2$ BRST Cohomology

In the Hilbert space one can form an inner product (*involution*)

$$\int d^4x \langle \Phi | \Phi \rangle := \int d^4x \int d^n\psi d\gamma d\beta dc \bar{\Phi}_{p,q} \bar{\psi}^{\mu_1} \wedge \cdots \wedge \bar{\psi}^{\mu_k} \wedge \underbrace{\delta^p(\gamma)}_{\bar{\beta}^p} \wedge \underbrace{\delta^q(\beta)}_{\bar{\gamma}^q} | \Phi \rangle$$

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Hence it holds:

**Vanishing Theorem** [Frenkel, Garland, Zuckerman]

$$H^\bullet(\mathcal{C}^{p,q}, Q) \begin{cases} \neq 0 & \text{for } H^0(\mathcal{C}^{p,p}, Q) \\ = 0 & \text{otherwise} \end{cases}$$

*The non-trivial cohomology is only in ghost degree zero.*

An alternative way to find the dimension of the cohomology groups relies on

$$\mathbb{P}(t) = \sum_k t^k \dim \hat{\mathcal{C}}_k^{\bullet, \bullet} \quad (2)$$

because of the **Euler-Poincaré principle**:

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**Observation:** for  $\mathcal{N} = 2$  spinning particle

$$\frac{1}{1-t} \mathbb{P}(t) = \dim H^{\bullet} \stackrel{\text{vanishing th.}}{\equiv} H^0(\hat{\mathcal{C}}^{p,p}, Q)$$

- We investigated also other **pictures**
- Topological sectors arise in the case of pseudoforms

# BV on the fields

Data for Batalin-Vilkovisky:  $T^*[-1]\mathcal{F}, \omega, \mathcal{Q}$ .

- The states in  $\hat{\mathcal{C}}^{p,q}$  are rearranged as  $\mathcal{F} = \underbrace{\mathcal{F}_0}_{\ni A_\mu(x), \varphi(x)} \oplus \underbrace{\mathcal{F}_1}_{\ni C(x)}$

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BRST of  $\mathcal{N} = 2$  spinning particle  $\Leftrightarrow$  classical, **non-interacting** BV of Yang–Mills field

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$$Q|\Phi\rangle =: \mathcal{Q}\Phi(x) \quad (3)$$

Look for the Hamiltonian vector fields to retain the action functional

$$\mathcal{Q} = \{S_{\text{free}} + (\Phi^* \mathcal{Q}\Phi)_{\text{free}}, -\} \quad (4)$$

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Observation:

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## A solution:

The action turns out to be:

$$\int d^4x \frac{1}{2} \langle \beta | Q_{\text{int}} Q Q_{\text{int}} | \beta \rangle + \frac{1}{3} \langle \beta | Q_{\text{int}} Q_{\text{int}} Q_{\text{int}} | \beta \rangle \equiv S[A, A^*, \varphi, \varphi^*, C, C^*]_{\text{BV YM}} \quad (5)$$

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- Perturbative quantization of the string field leads to BV,  $L_\infty$  algebra

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- Perturbative quantization of the string field leads to BV,  $L_\infty$  algebra
- A spinning particle on superworldline lends a BV structure to the space of fields in the target
- $\mathcal{N} = 2$  spinning particle  $\Leftrightarrow$  open string
- the full action functional of the target Yang–Mills theory is retrieved

