

# Asymptotically Kerr-de Sitter Spacetimes: Necessary Conditions

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(Unpublished work with Jaroslaw Kopinski)

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**Solution:** Conformal Killing–Yano 2-forms

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**Study obstructions asymptotically.**

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*We can study asymptotically Kerr-NUT-dS spacetimes by studying “conformally-invariant” series expansion of  $Q^{\Omega^2 g}$  near  $\partial\bar{M}$  in  $\Omega$ .*

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$$0 \stackrel{\partial\bar{M}}{=} \top(\nabla_n - H)Q_{ab} + \bar{\nabla}_{[a} Q_{nb]}^\top$$

$$0 \stackrel{\partial\bar{M}}{=} (d-2)n^b \bar{g}_a^c (\nabla_n Q)_{bc} - (d-2)H Q_{na}^\top + \bar{\nabla}^b Q_{ab}^\top$$

$$0 \stackrel{\partial\bar{M}}{=} \overline{CKY}(Q^\top)_{abc}$$

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and “normal vector” like section  $[\rho := -\frac{1}{d}(\Delta\Omega + J\Omega)]$

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Conformally-invariant “normal derivative”:

$$I \cdot \hat{D} := (\nabla_n + w\rho) - \frac{\Omega}{d+2w-2}(\Delta + wJ).$$

# Constructing higher-order conformally-invariant constraints II

**Step 1:** “Insert”  $\mathcal{Q}_{abc}$  into some tractor bundle, so  
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2 evolution- + 2 boundary constraint-type equations. This is the expected general trend.

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**Boundary constraint-type equations characterize “how” Kerr-de Sitter a spacetime can be.**

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If  $\mathcal{Q} = \mathcal{O}(\Omega^3)$ , then:

$$(\nabla_n^3 Q_{ab})^\top + \text{more} \stackrel{\partial \bar{M}}{=} 0$$

$$n^a \bar{g}_c^b \nabla_n^3 Q_{ab} + \text{more} \stackrel{\partial \bar{M}}{=} 0$$

$$\mathring{IV}_{(a} Q_{b)c}^\top \stackrel{\partial \bar{M}}{=} 0$$

$$\mathring{IV}_{a[b} Q_{nc]}^\top - \frac{1}{2} \bar{g}_{a[b} \mathring{IV}_{c]} \cdot Q_n^\top \stackrel{\partial \bar{M}}{=} 0$$

As  $\mathring{IV}$  is the undetermined Neumann data, these constraints are *interesting*.

**Boundary constraint-type equations characterize “how” Kerr-de Sitter a spacetime can be. Matter!**

Thank you

Sam Blitz

Thank you!